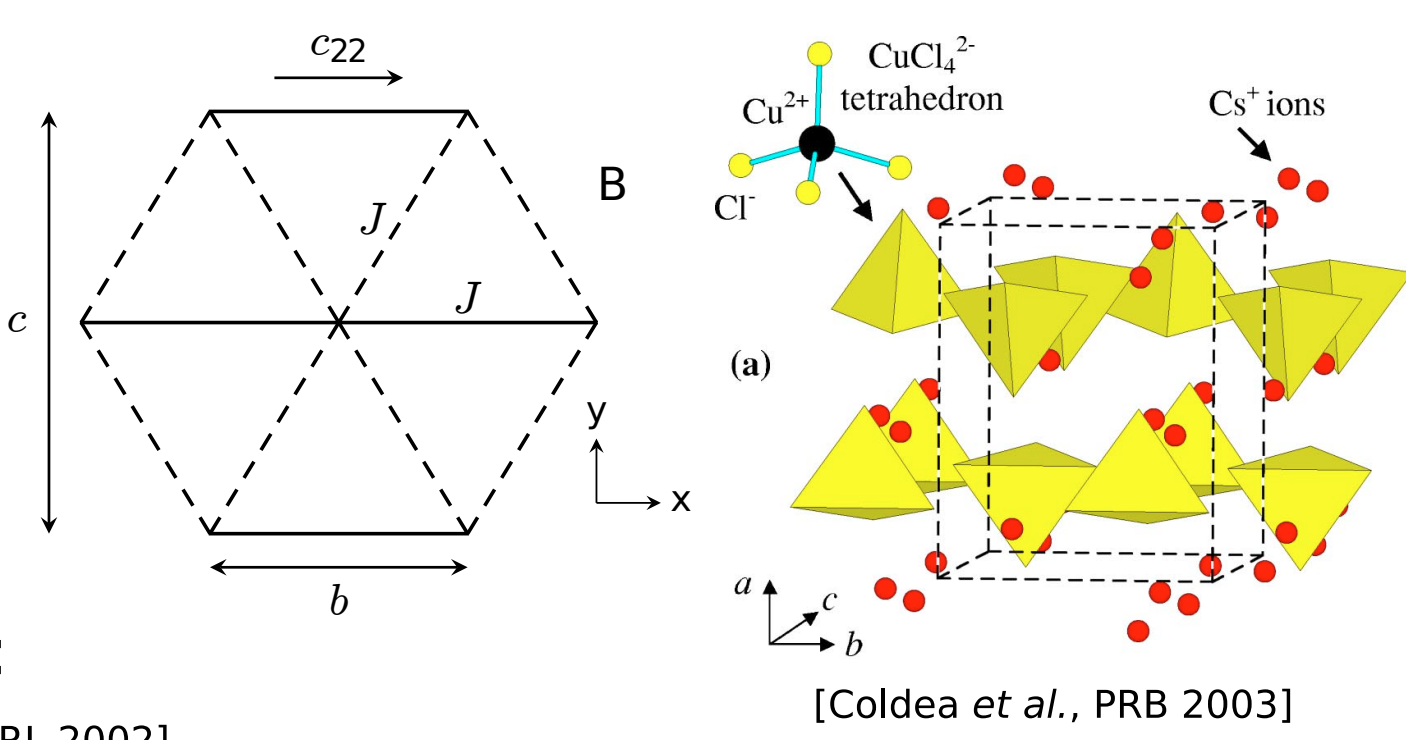


### Cs<sub>2</sub>CuCl<sub>4</sub>

- motivation: ultrasound physics ( $c_{22}$ -mode) in the spin-liquid phase of Cs<sub>2</sub>CuCl<sub>4</sub>
- crystal structure: orthorhombic ( $Pnma$ )
- Cu<sup>2+</sup>: carry spin 1/2 (orbital quenching) and form a triangular lattice in each plane
- effective 2D Hamiltonian with  $B \parallel a$ :  

$$\mathcal{H} = \frac{1}{2} \sum_{ij} J_{ij} S_i \cdot S_j - h \sum_i S_i^z$$
with  $h = g\mu_B B$ ,  $g = 2.2$
- two nearest neighbor exchange couplings:  
 $J = 4.34\text{K}$  and  $J' = 1.49\text{K}$  ( $k_B \equiv 1$ ) [Coldea et al., PRL 2002]  
 (from neutron scattering data in the ferromagnetic phase)



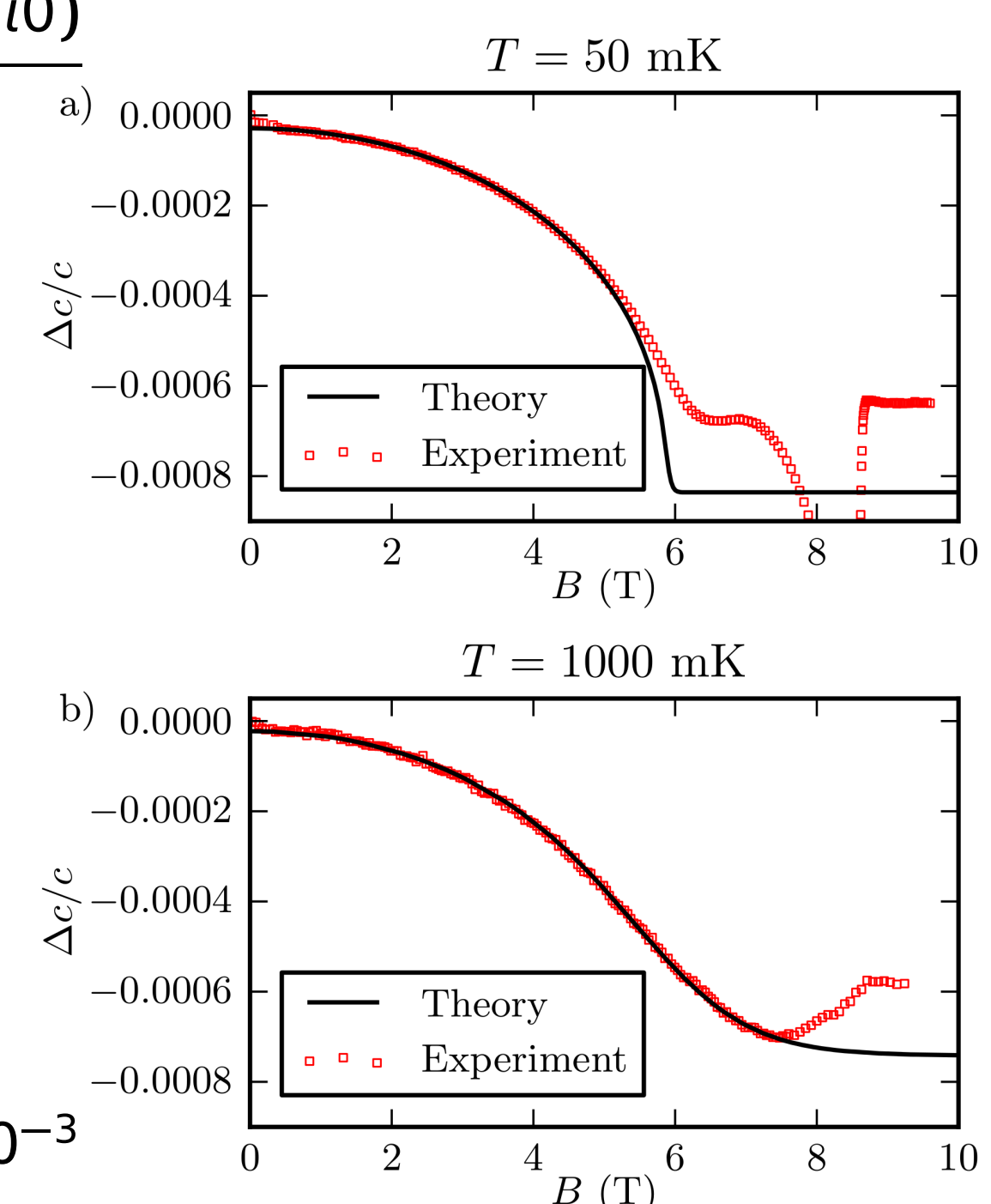
### Sound velocity ( $c_{22}$ -mode)

- renormalized dispersion:  $\tilde{\omega}_q = \omega_q + \frac{\text{Re}\Pi(q, \omega_q + i0)}{2\omega_q}$
- renormalized sound velocity:  $\tilde{c} = \lim_{q \rightarrow 0} \tilde{\omega}_q/q$
- shift of the sound velocity:  $\Delta c = \tilde{c} - c$   

$$\Delta c/c = g_1 c^{(1)} + g_2 c^{(2)}$$

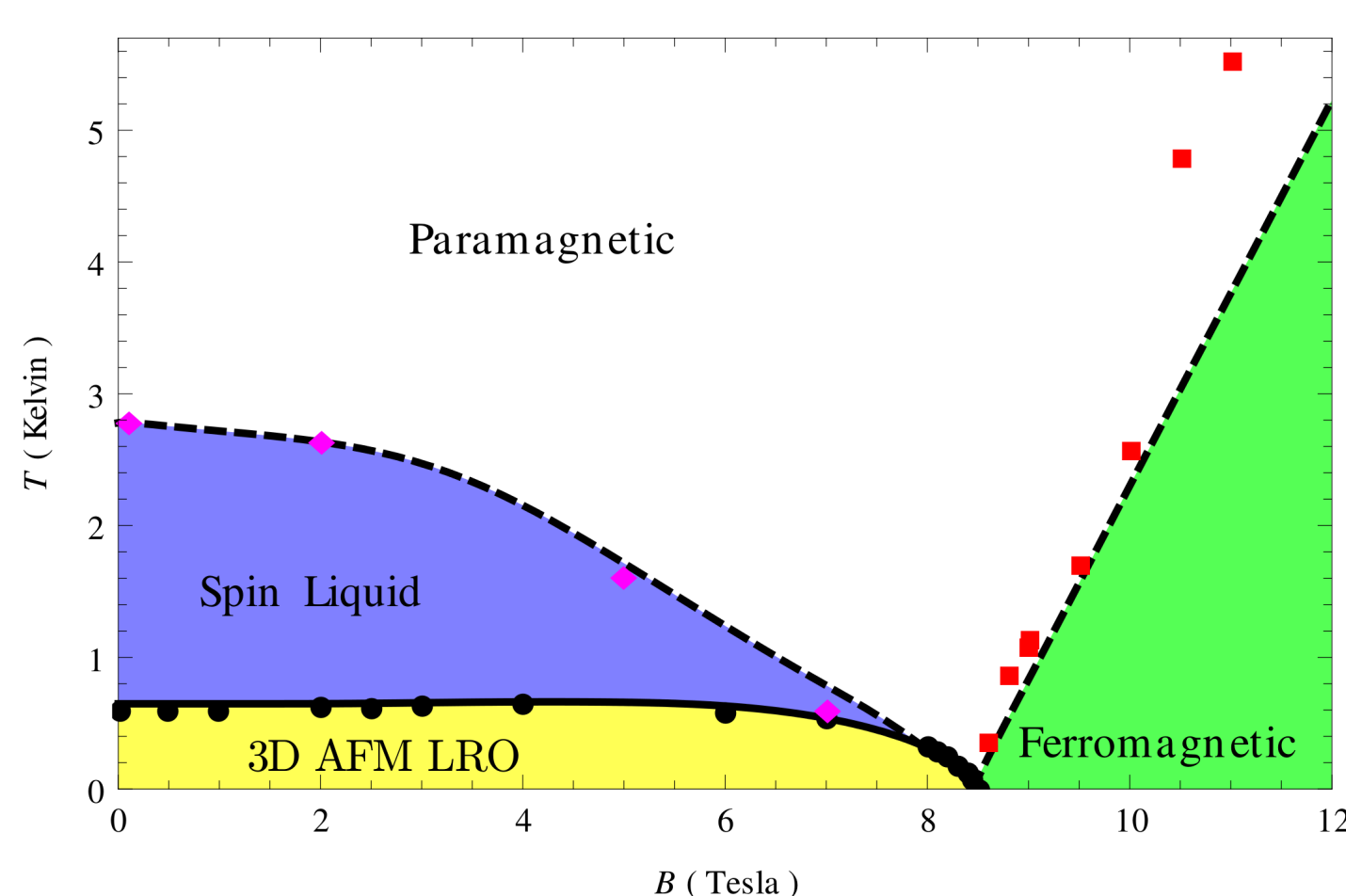
$$c^{(1)} = \int_{-\pi}^{\pi} \frac{dk}{2\pi} \frac{J f'(\xi_k)}{v_k - c} \frac{v_k}{v_k - c} [2s - Z \cos k]^2$$

$$c^{(2)} = s^2 - Z^2/4$$
- dimensionless coupling constants:  
 $g_1 = [J^{(1)}b]^2/(2Mc^2)$ ,  $g_2 = J^{(2)}b^2/(2Mc^2)$
- group velocity of the fermions:  $v_k = ZJb \sin k$
- Cs<sub>2</sub>CuCl<sub>4</sub>:  $c/(Jb) \approx 6.8 \Rightarrow c^{(2)} \gg c^{(1)}$  ( $\hbar \equiv 1$ )
- fit to experimental data:  $g_1 \approx 0$ ,  $g_2 \approx -1.1 \times 10^{-3}$



### Phase diagram: spin-liquid phase

- spin-liquid phase: no long-range order, but strong short-range correlations
- inelastic neutron scattering: highly dispersive excitation continua, signature of spin-1/2 spinons [Coldea et al., PRB 2003]
- quasi-one-dimensional behavior due to frustration and quantum fluctuations
- dilute Bose gas quantum critical point at  $B_c = 8.5\text{T}$ : BEC of magnons [Radu et al., PRL 2005]

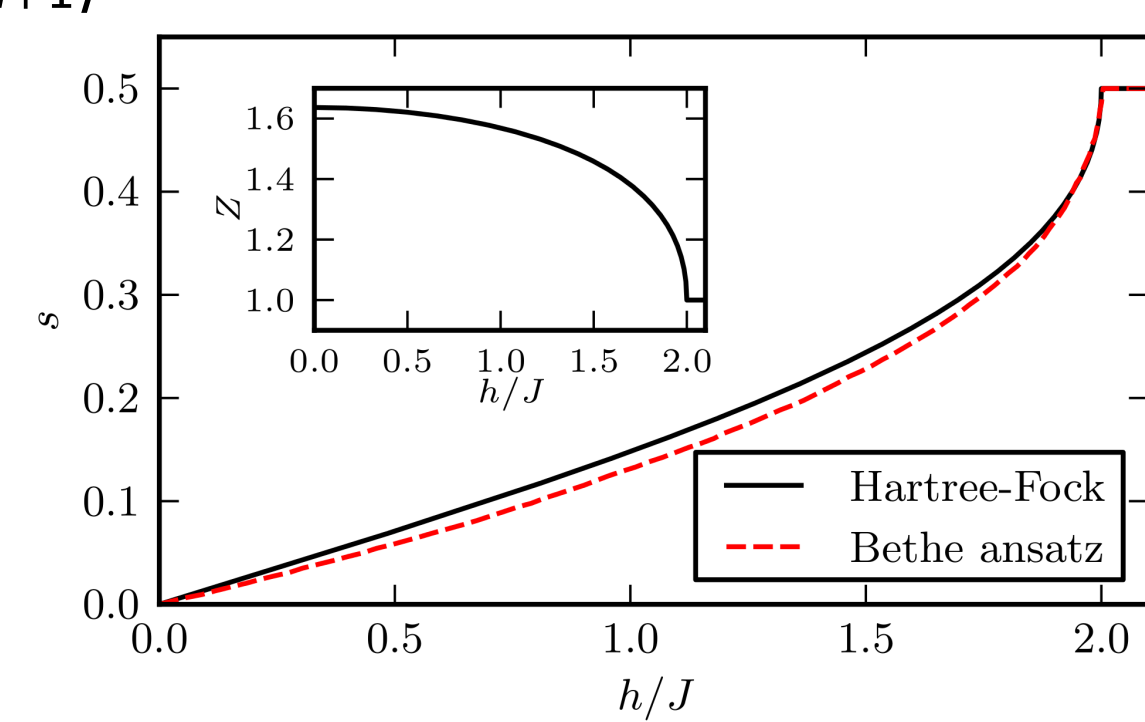


### Heisenberg chain coupled to phonons

- spin-phonon Hamiltonian:  $\mathcal{H} = \sum_n J_n [S_n \cdot S_{n+1} - 1/4] - h \sum_n S_n^z + \mathcal{H}_2^p$
- one-dimensional phonons:  $\mathcal{H}_2^p = \sum_q \left[ \frac{P_{-q} P_q}{2M} + \frac{M}{2} \omega_q^2 X_{-q} X_q \right]$ ,  $\omega_q = c \frac{|q|}{b}$
- exchange-striction mechanism:  $J_n \approx J + J^{(1)}(X_{n+1} - X_n) + \frac{J^{(2)}}{2}(X_{n+1} - X_n)^2$
- Jordan-Wigner transformation to spinless fermions:  
 $S_n^+ = c_n^\dagger (-1)^n e^{i\pi \sum_{j < n} c_j^\dagger c_j}$ ,  $S_n^- = (-1)^n e^{-i\pi \sum_{j < n} c_j^\dagger c_j} c_n$ ,  $S_n^z = c_n^\dagger c_n - 1/2$   

$$\mathcal{H} = -\frac{1}{2} \sum_n J_n (c_n^\dagger c_{n+1} + c_{n+1}^\dagger c_n + c_n^\dagger c_n + c_{n+1}^\dagger c_{n+1}) + \sum_n J_n c_n^\dagger c_n c_{n+1}^\dagger c_{n+1} - h \sum_n c_n^\dagger c_n + N h/2 + \mathcal{H}$$
- self-consistent Hartree-Fock theory:  

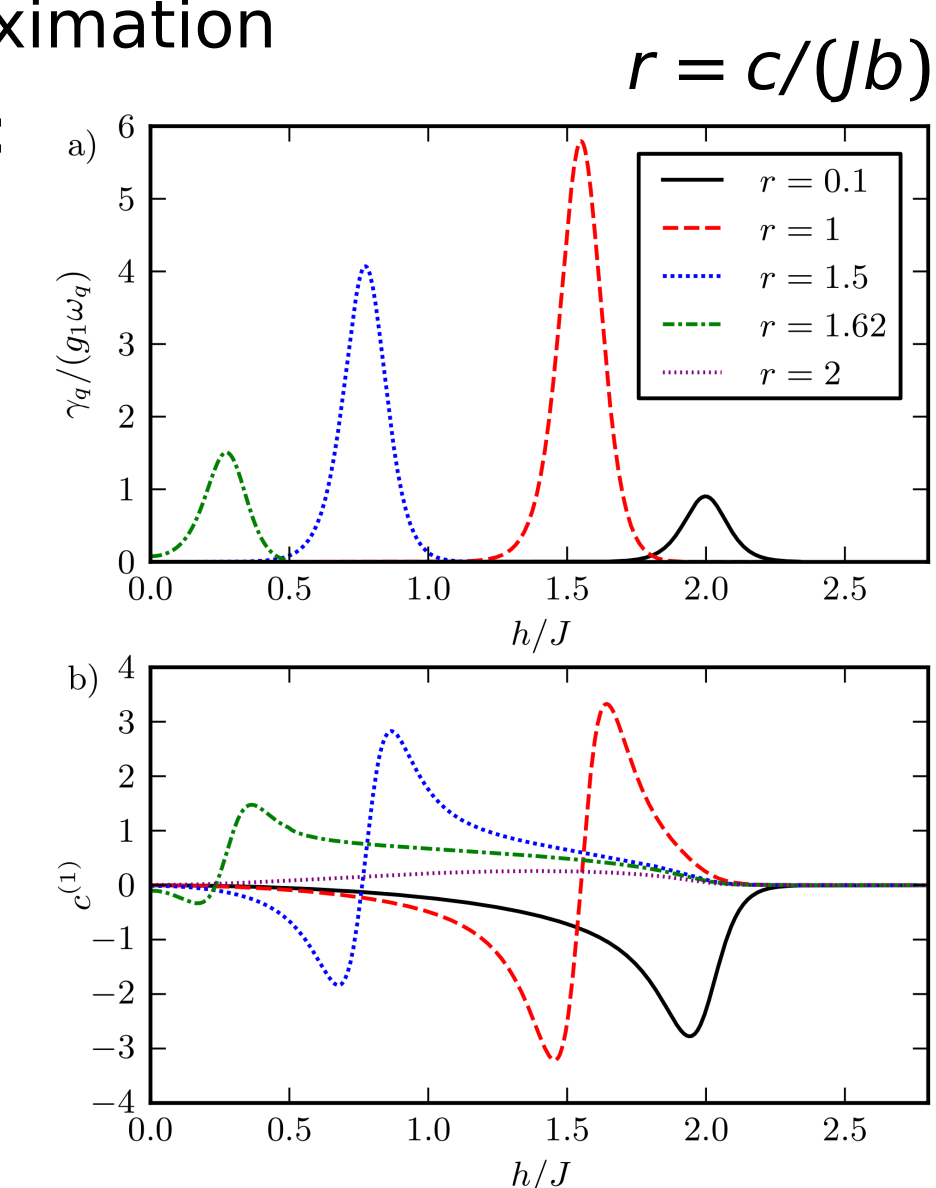
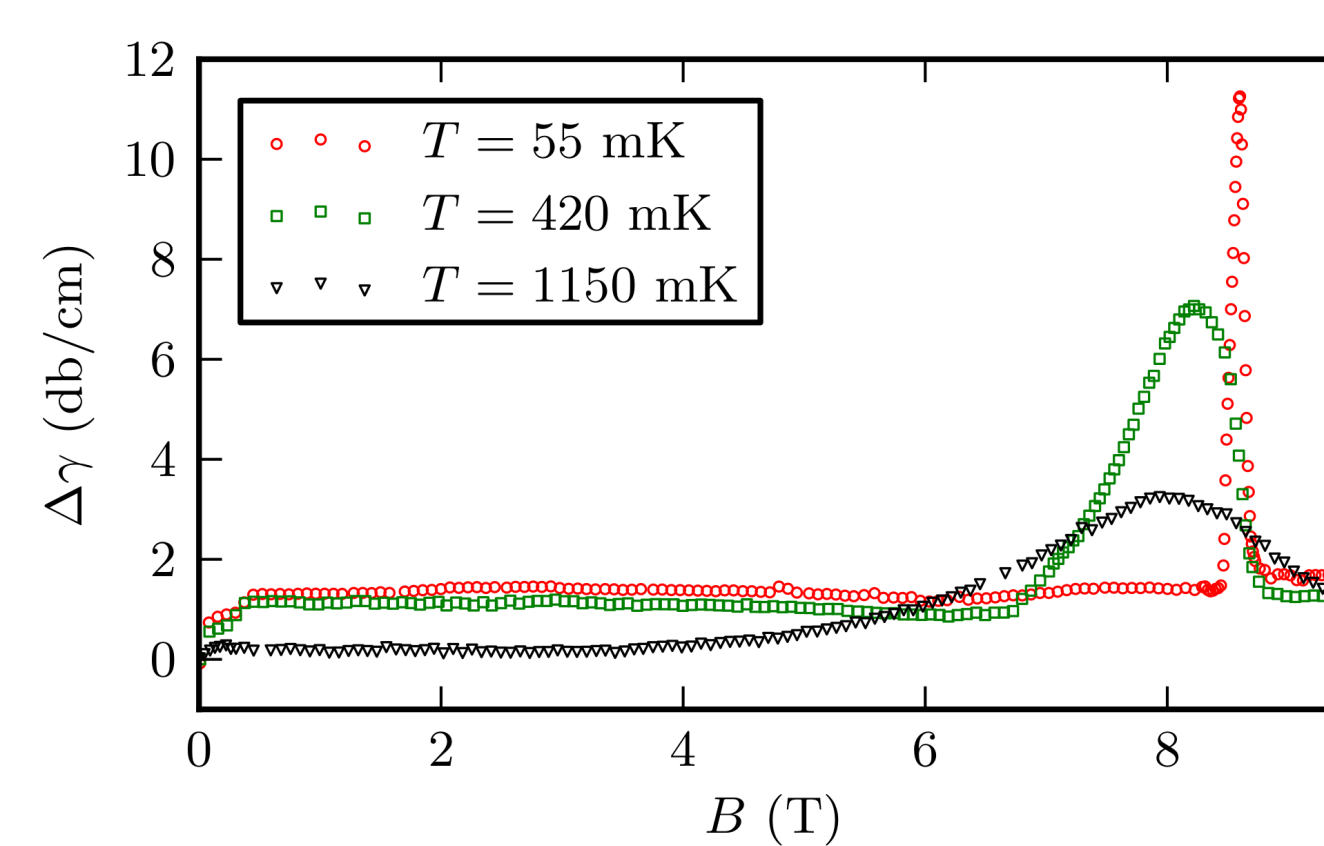
$$c_n^\dagger c_n c_{n+1}^\dagger c_{n+1} \approx \rho (c_{n+1}^\dagger c_{n+1} + c_n^\dagger c_n) - \rho^2 - \tau (c_n^\dagger c_{n+1} + c_{n+1}^\dagger c_n) + \tau^2$$
- self-consistency conditions:  $\rho = \langle c_n^\dagger c_n \rangle$ ,  $\tau = \langle c_n^\dagger c_{n+1} \rangle$
- fermion dispersion:  $\xi_k = -ZJ \cos k + 2s_j - h$
- renormalization of the hopping:  $Z = 1 + 2\tau$
- dimensionless magnetization:  $s = \rho - 1/2$



### Ultrasound attenuation ( $c_{22}$ -mode)

- phonon damping:  

$$\gamma_q = -\frac{\text{Im}\Pi(q, \omega_q + i0)}{2\omega_q} = \frac{\pi}{2M\omega_q} \int_{-\pi}^{\pi} \frac{dk}{2\pi} (f_k - f_{k+q}) |\Gamma_3(k, q)|^2 \delta(\xi_k - \xi_{k+q} + \omega_q)$$
- limit  $q \rightarrow 0$ :  $\frac{\gamma_q}{\omega_q} \sim \frac{g_1 c J \Theta(v_*^2 - c^2)}{2v_* \sqrt{1 - c^2/v_*^2}} [-f'(\xi_+) V_+^2 - f'(\xi_-) V_-^2]$ ,  $V_\pm = 2s \pm Z \sqrt{1 - c^2/v_*^2}$ ,  $\xi_\pm = J V_\pm - h$
- only finite if  $v_* = ZJb > c \Rightarrow$  no damping in our approximation
- natural explanation for the deviations at larger fields: dilute Bose gas quantum critical point at  $B_c = 8.5\text{T}$



### Spin-phonon coupling

- approximate spin-phonon Hamiltonian:  $\mathcal{H} = F_0 + \sum_k \xi_k c_k^\dagger c_k + \mathcal{H}_2^p + \delta\mathcal{H}_2^p + \mathcal{H}_3^{SP} + \mathcal{H}_4^{SP}$   

$$F_0/N = h/2 + J(\tau^2 - \rho^2)$$
,  $\delta\mathcal{H}_2^p = 2J^{(2)}(\tau^2 - \rho^2) \sum_q \sin^2(q/2) X_{-q} X_q$   

$$\mathcal{H}_3^{SP} = \frac{1}{\sqrt{N}} \sum_{k'kq} \delta_{k',k+q}^* \Gamma_3(k, q) c_{k'}^\dagger c_k X_q$$
  

$$\mathcal{H}_4^{SP} = \frac{1}{2N} \sum_{k'kq_1q_2} \delta_{k',k+q_1+q_2}^* \Gamma_4(k, q_1, q_2) c_{k'}^\dagger c_k X_{q_1} X_{q_2}$$
  

$$\Gamma_3(k, q) \approx -iqJ^{(1)}[Z \cos k - 2s]$$
,  $\Gamma_4(k, q_1, q_2) \approx q_1 q_2 J^{(2)}[Z \cos k - 2s]$
- propagator of the phonon field  $X_q$ :  $\langle X_{q,\omega} X_{-q,-\omega} \rangle = \frac{\beta N}{M} \frac{1}{\omega^2 + \omega_q^2 + \Pi(q, i\omega)}$
- phonon self-energy:  $\Pi(q, i\omega) = \Pi_2(q) + \Pi_3(q, i\omega) + \Pi_4(q)$   

$$\Pi_2(q) = [J^{(2)}/M][\tau^2 - \rho^2] 4 \sin^2(q/2)$$
,  $\Pi_4(q) = \frac{1}{MN} \sum_k f_k \Gamma_4(k, q, -q)$   

$$\Pi_3(q, i\omega) = \frac{1}{MN} \sum_k \frac{f_k - f_{k+q}}{\xi_k - \xi_{k+q} + i\omega} |\Gamma_3(k, q)|^2$$
,  $f_k = \frac{1}{e^{\beta \xi_k} + 1}$

### Research plan

- calculate the correlation functions in the vicinity of the dilute Bose gas quantum critical point by combining the hard-core boson formalism with the functional renormalization group, going beyond the usually used ladder approximation

$$\mathcal{H}_U = \sum_k (\epsilon_k - \mu) b_k^\dagger b_k + \frac{1}{2N} \sum_{q,k,k'} (U_q + U) b_{k+q}^\dagger b_{k'-q}^\dagger b_{k'} b_k$$

$$\epsilon_k = J_k^D - J_{\min}^D, \quad \mu = \frac{J_0^D - J_{\min}^D}{2} - h = h_c - h, \quad J_k^D = J_k + D_k$$

$$J_k = \sum_R J(R) \cos(k \cdot R), \quad D_k = \sum_R D(R) \sin(k \cdot R)$$

- generalize the hard-core boson formalism to take the coupling of hard-core bosons to the lattice vibrations (phonons) into account
- study ultrasound in the vicinity of the dilute Bose gas quantum critical point and compare with the experiments of B1
- study driven dissipative dilute Bose gas, non-equilibrium time evolution
- connections to projects **B1 B2 B3 B4**

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